

Velocity-dependent transverse momentum distribution of projectile- like fragments at 95 MeV/u

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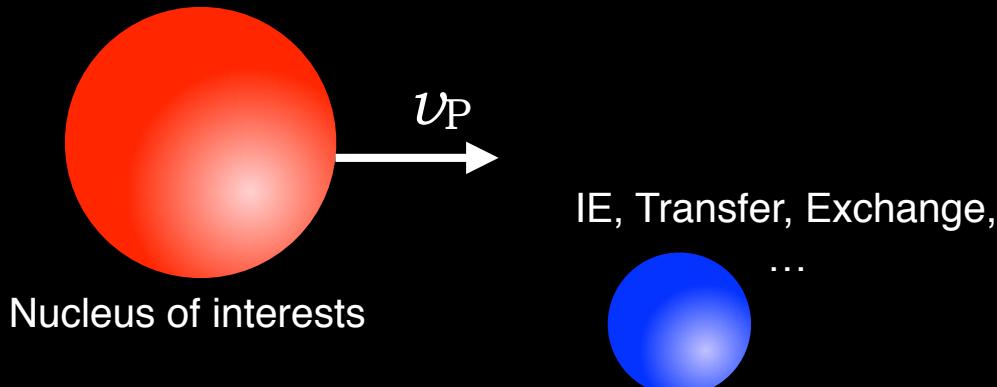
r280n collaboration

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Introduction

Investigation of unstable nuclei

Inverse kinematics scheme



Preparation of radioactive nuclear beams (RNBs)

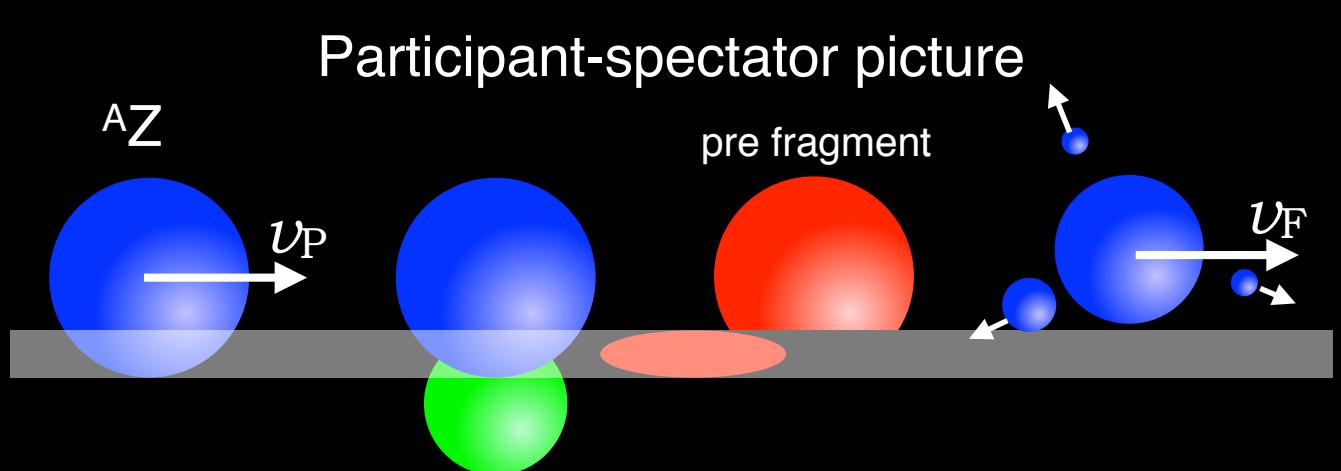
- 1) Wide range in nuclear chart
- 2) High quality as a beam

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Fragmentation process



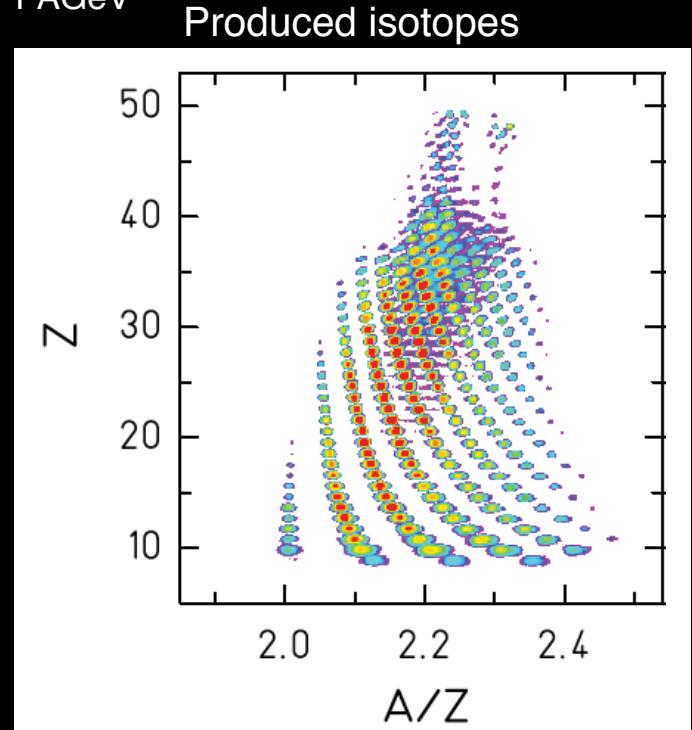
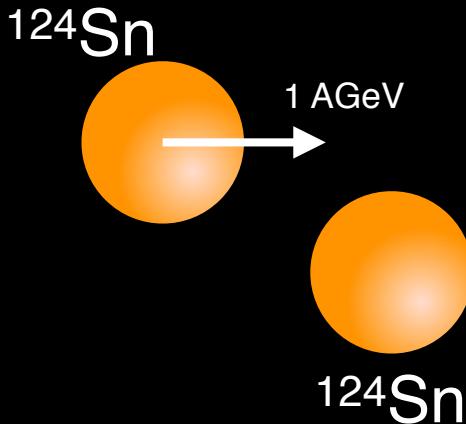
Satisfies two requirements

- 1) High productivity for wide range of isotopes
- 2) High beam quality

→ Usually applied at RNB facilities

Production of a wide range of isotopes

Fragments prod. from $^{124}\text{Sn} + ^{124}\text{Sn}$ at 1 AGeV
V. Föhr et al., Phys. Rev. C 84, 054605 (2011).



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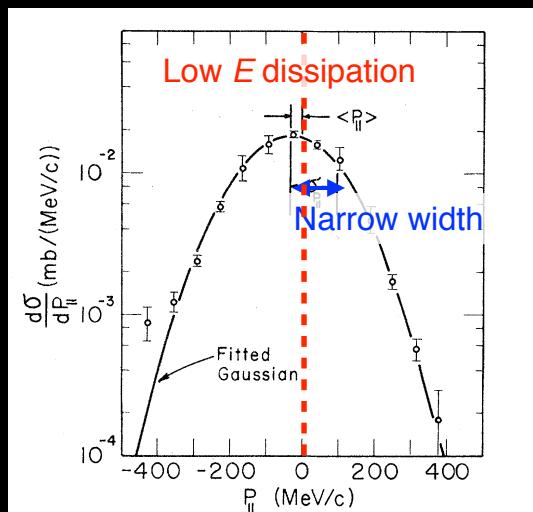
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High quality as secondary beam

Momentum distribution of fragments

- E dissipation : low
- Narrow width
- Simple & well studied at relativistic E

P_L distribution of ^{10}Be produced from ^{12}C (2.1 GeV/nucleon)+ ^9Be
D.E. Greiner et al., Phys. Rev. Lett. 35(1975) 152



$$P_0=2880 \text{ MeV/c/A}$$

Therefore

- Separation of objective isotope with high efficiency

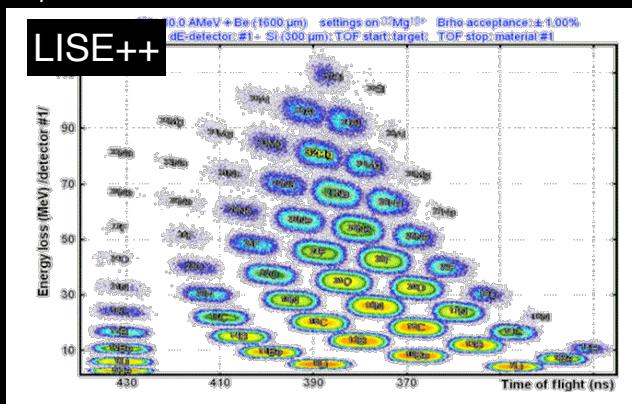
Formulation for practical use

Transportation/separation of fragments through fragment separator

Performance of separator is simulated by means of

LISE++, MOCADI, ...

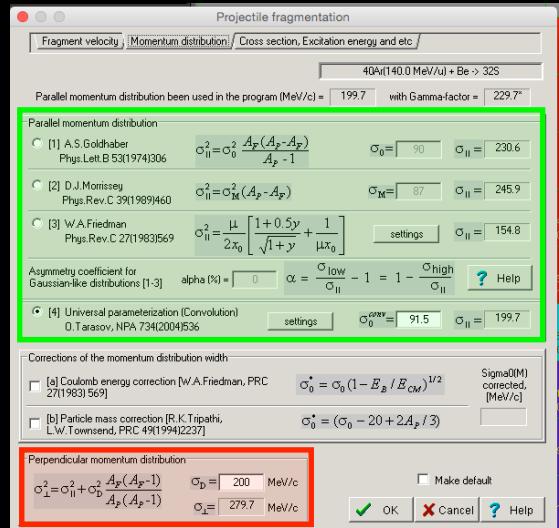
Ex. ^{40}Ca (80 AMeV) + Be optimized for ^{32}Mg
<http://lise.nscl.msu.edu/introduction.html>



Key parameters for simulation

$\sigma_{\text{prod.}}$

P_L , P_T distribution



Few systematic studies on P_T

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Object of this talk

Remarkable and systematic correlation between P_T distribution and P_L of fragmentation products at $E \sim 100$ MeV/u

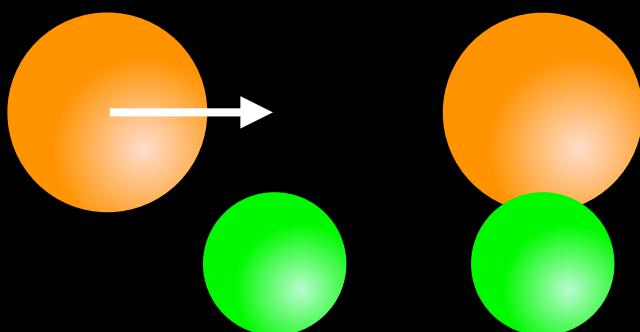
1. Earlier works on P_L , P_T distributions
2. Experimental (RIPS-RIKEN)
3. Correlation obtained from experimental results
4. Comparison with microscopic dynamic model
5. Conclusions

Previous studies on momentum distributions

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Isotropic distribution at relativistic E

$E \gtrsim 1 \text{ GeV/u}$

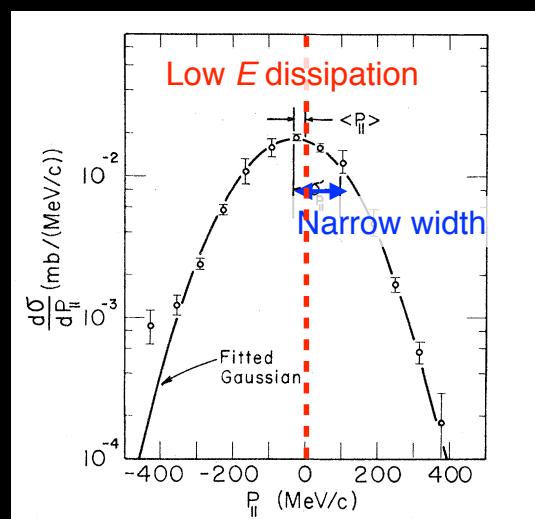


- Gaussian-type distribution with small energy dissipation
- Isotropic distribution

$\sigma_L = \sigma_T$ to an accuracy of 10%

→ Contribution of Fermi momentum

P_L distribution of ^{10}Be
 $^{12}\text{Ca}(2.1 \text{ GeV/nucleon}) + ^9\text{Be}$
D.E. Greiner et al., Phys. Rev. Lett. 35(1975) 152

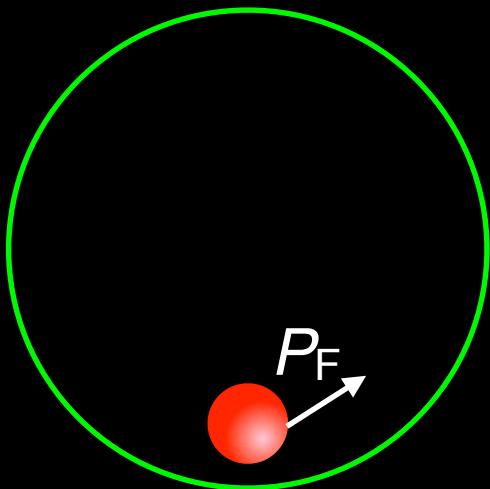


Model based on Fermi momentum

Assumption :
Independent removal of
nucleons in projectile



Momentum distribution of fragments
corresponds to statistical sum of
Fermi momentum for each removed
nucleon.



P_F : Fermi momentum of
nucleon in projectile

Formulation proposed by Goldhaber

$$\sigma_{\text{GH}}^2 = \frac{A_F(A_P - A_F)}{A_P - 1} \sigma_0^2$$

$$\sigma_0 = \frac{P_F}{\sqrt{5}} \sim 100 \text{ MeV/c}$$

A.s. Goldhaber, Phys. Lett. B 53, 306 (1974).

Simple and successfully applied to
a wide range of reaction system

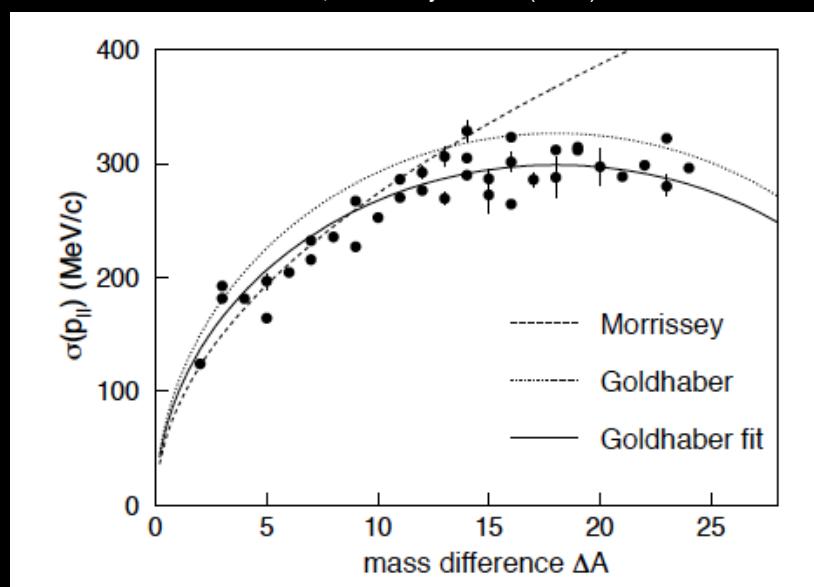
II

Success of Goldhaber model

Width of P_L distribution of fragments

$^{36}\text{Ar}(1.05 \text{ GeV/nucleon}) + ^9\text{Be}$

M. Caamano et al., Nucl. Phys. A 733(2004) 187.

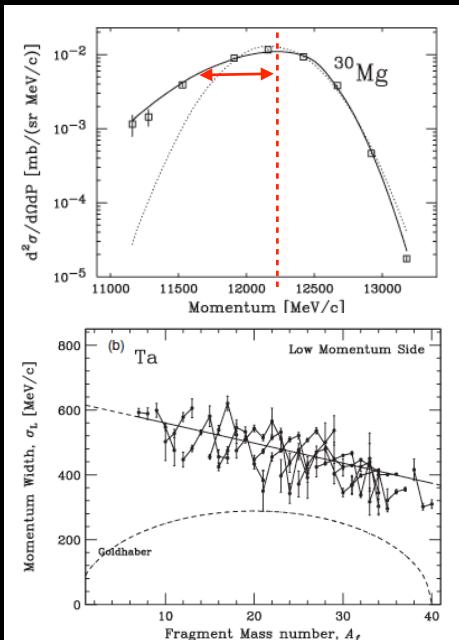


$$\sigma_0 = 98.2 \text{ MeV/c} \quad \frac{P_F}{\sqrt{5}} \sim 100 \text{ MeV/c}$$

Deviation from isotropic dist. at $E \sim 100$ MeV/u

P_L distribution : Low momentum tail

P_L distribution at 90 MeV/u
 $^{40}\text{Ar} + ^9\text{Be} \rightarrow ^{30}\text{Mg} + X$



M. Notani et al., Phys. Rev. C 76 (2007) 044605.

Universal parametrization obtained from experimental results

O. Tarasov, Nucl. Phys. A 734 (2004) 536.

$$f(P_L) = \exp\left(\frac{P_L}{\tau}\right) \cdot \left[1 - \operatorname{erf}\left(\frac{P_L - P_0 + \sigma_{\text{pf}}^2 / \tau - s \cdot \tau}{\sqrt{2\sigma_{\text{pf}}}}\right) \right]$$

$$\tau = \text{coeff} \cdot \frac{\sqrt{A_F \cdot E_S}}{\beta}$$

$$\sigma_{\text{pf}} = \beta \sigma_{\text{pf}}^2 \frac{A_F(A_P - A_F)}{A_P - 1}$$

Formulated momentum distributions have been incorporated into simulation.

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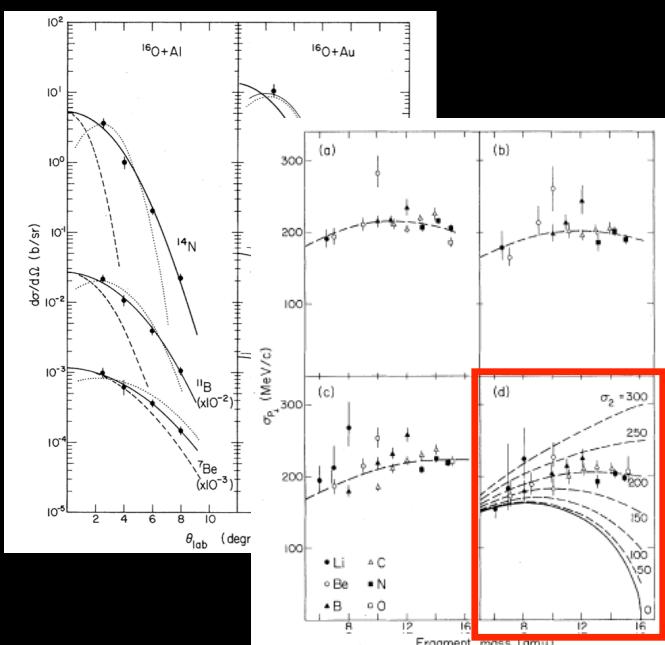
Deviation from isotropic dist. at $E \sim 100$ MeV/u

P_T distribution : Additional width

P_T distribution observed at 90 MeV/u



K van Bibber et al., Phys. Rev. Lett. 43 (1979) 840.



Additional width due to orbital deflection by target nucleus

$$\sigma_T^2 = \sigma_{\text{GH}}^2 + \sigma_D^2$$

Empirical formulation for σ_D

$$\sigma_D^2 = \frac{A_F(A_F - 1)}{A_P(A_P - 1)} \sigma_2^2$$

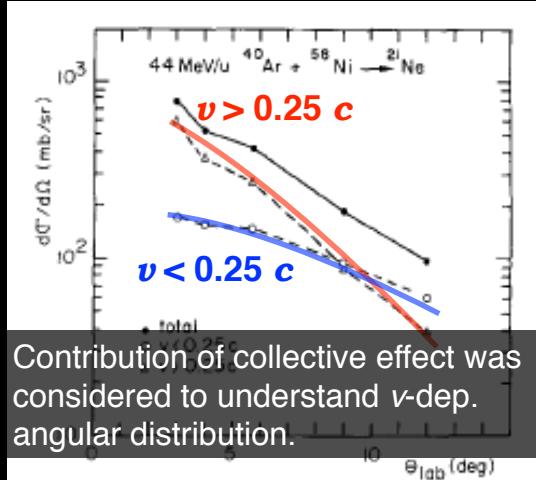
Usually applied with $\sigma_2 \sim 200$ MeV/c.

However, the reliability of the formulation is doubtful

- scarce systematic measurements at $E \sim 100$ MeV/u
- no measurements as a func. of v . suggested from asymmetric P_L dist.

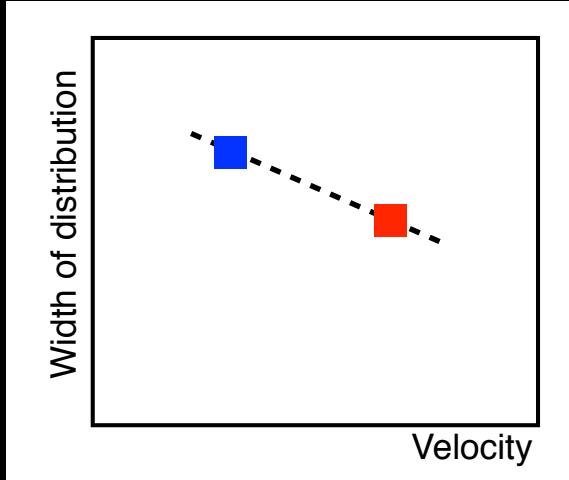
Velocity dependent P_T distribution

Angular dist. observed at 44 MeV/u
 $^{40}\text{Ar} + ^{58}\text{Ni} \rightarrow ^{21}\text{Ne} + X$



V. Borrel et al., Z. Phys. A 314 (1982) 191.

Expected trend for width of P_T distribution



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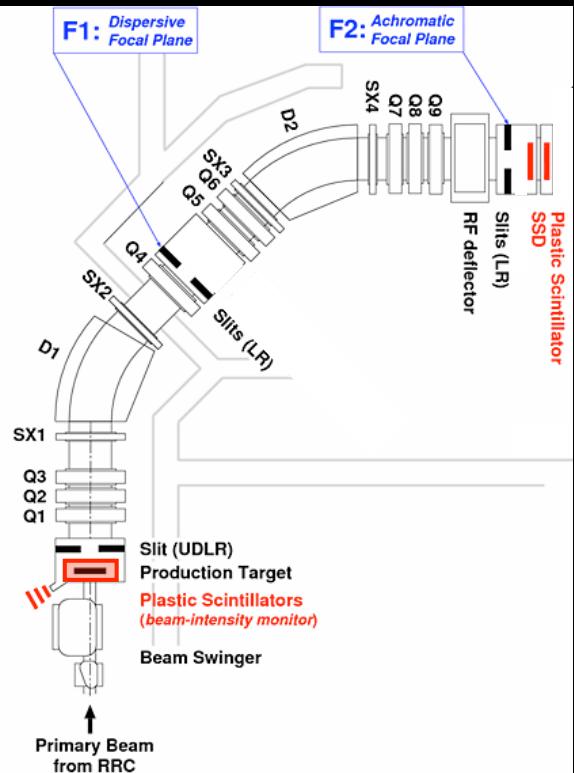
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Experimental

Experimental : fragment separator

RIKEN Ring cyclotron + RIPS



Reaction : ^{40}Ar (95 MeV/u) + ^9Be

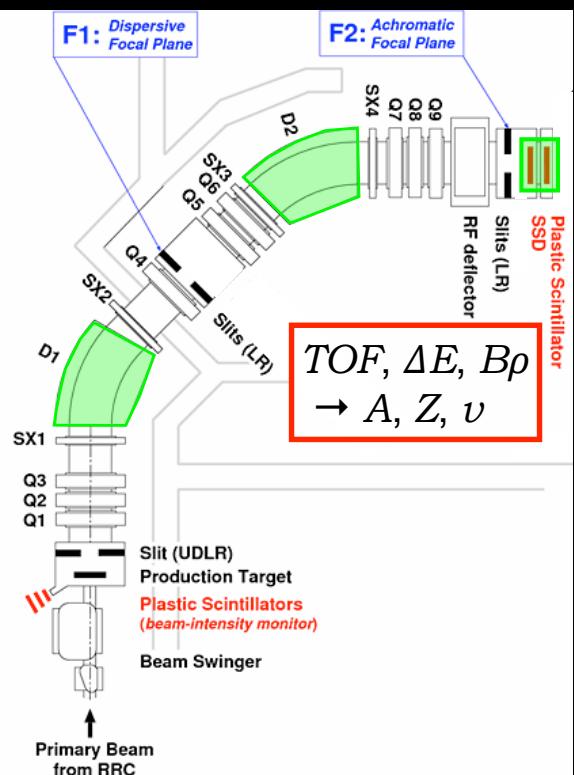
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Experimental : fragment separator

RIKEN Ring cyclotron + RIPS



Reaction : ^{40}Ar (95 MeV/u) + ^9Be

Identification : TOF, ΔE , $B\rho$

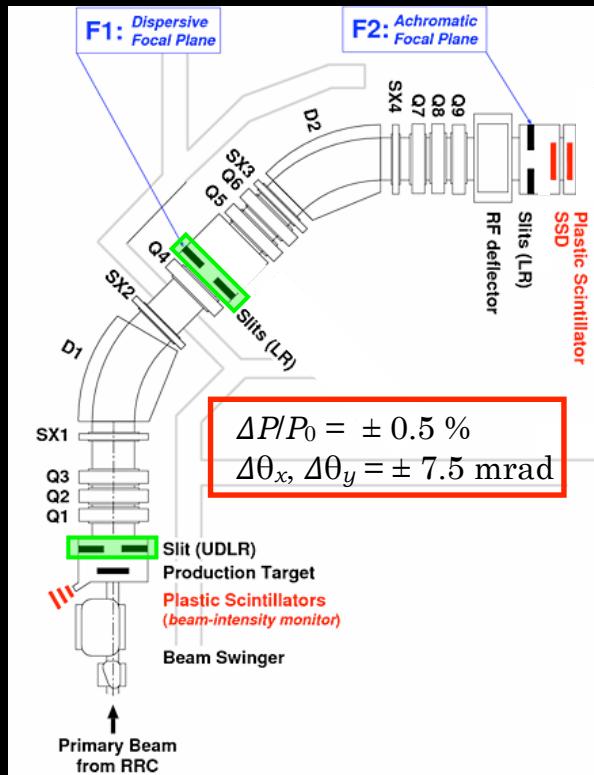
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Experimental : fragment separator

RIKEN Ring cyclotron + RIPS



Reaction : ^{40}Ar (95 MeV/u) + ^9Be

Identification : TOF, ΔE , $B\rho$

Acceptance : Slit after target, F1

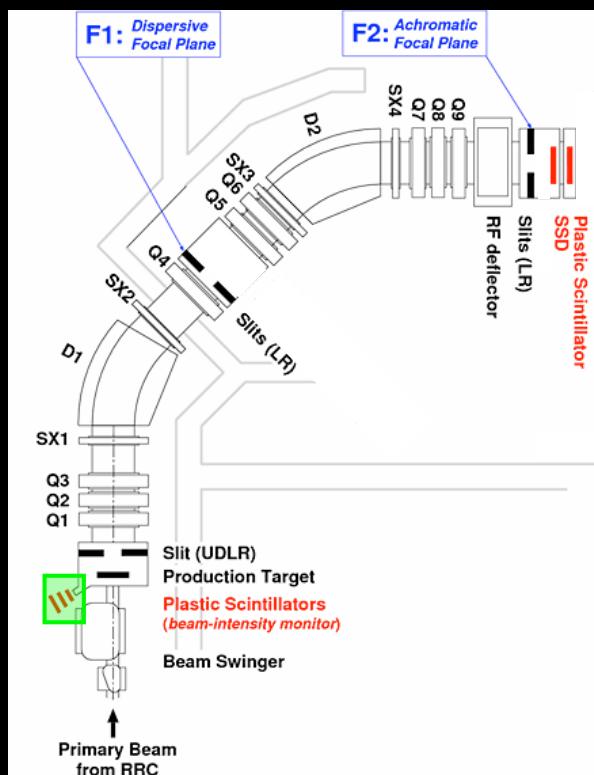
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Experimental : fragment separator

RIKEN Ring cyclotron + RIPS



Reaction : ^{40}Ar (95 MeV/u) + ^9Be

Identification : TOF, ΔE , $B\rho$

Acceptance : Slit after target, F1

Beam intensity monitor : PL@target
Ambiguity of beam intensity : $\sim \pm 5\%$

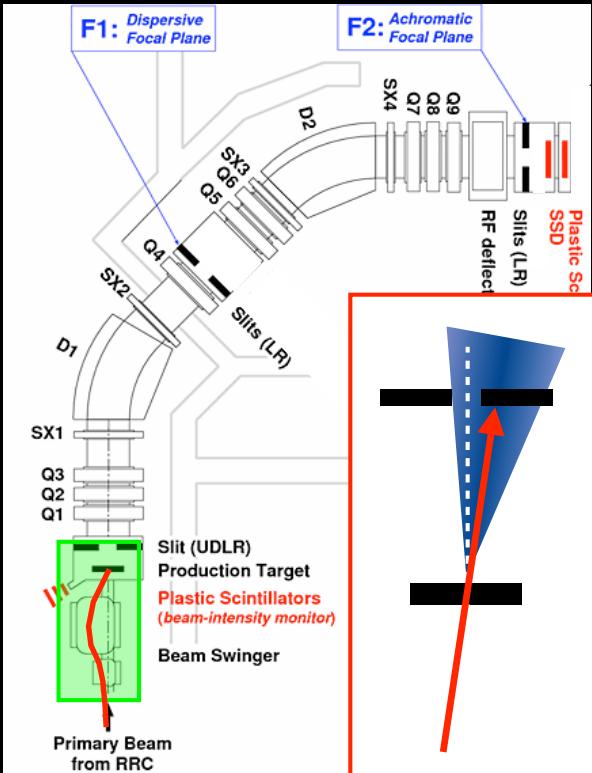
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Experimental : fragment separator

RIKEN Ring cyclotron + RIPS



Reaction : ^{40}Ar (95 MeV/u) + ^9Be

Identification : TOF, ΔE , $B\rho$

Acceptance : Slit after target, F1

Beam intensity monitor : PL@target

Ambiguity of beam intensity : $\sim \pm 5\%$

Angular distribution :
Beam swinger + slit after target

Keep optical axis of RIPS at any angle setting
→ Constant values for transmission and detection efficiencies of fragmentation products.

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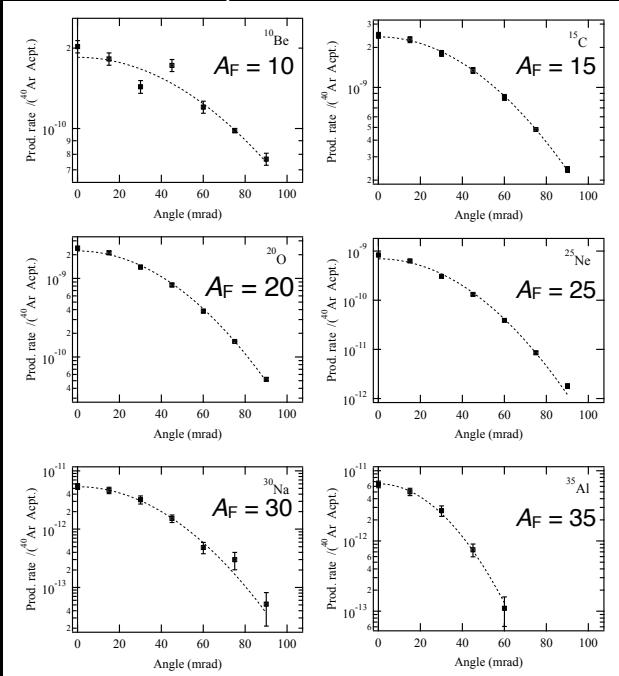
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Results & Analysis

Width of P_T distribution

Obtained from angular distribution

$$B\rho = 3.600 \text{ Tm}$$



S. Momota et al., Phys. Rev. C 92, 024608 (2015).

Fitting with a Gaussian function

$$N(\theta) = A \exp \left\{ -\frac{\theta^2}{2\sigma_\theta^2} \right\}$$

with considering

- 1) Finite angular acceptance $\pm 7.5 \text{ mrad}$
- 2) Angular straggling in target evaluated by ATIMA
- 3) Emittance of primary beam assumed to be neglected

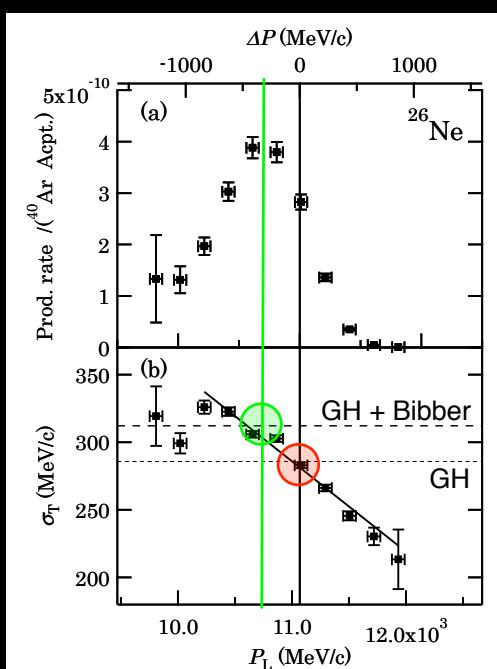
Width of P_T distribution

$$\sigma_T = P_L \times \sigma_\theta$$

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Correlation between σ_T and velocity



ΔP : shift from primary beam velocity

P_L distribution

Deceleration : $\sim 300 \text{ MeV/c}$

Larger width for low P_L

Remarkable decreasing trend

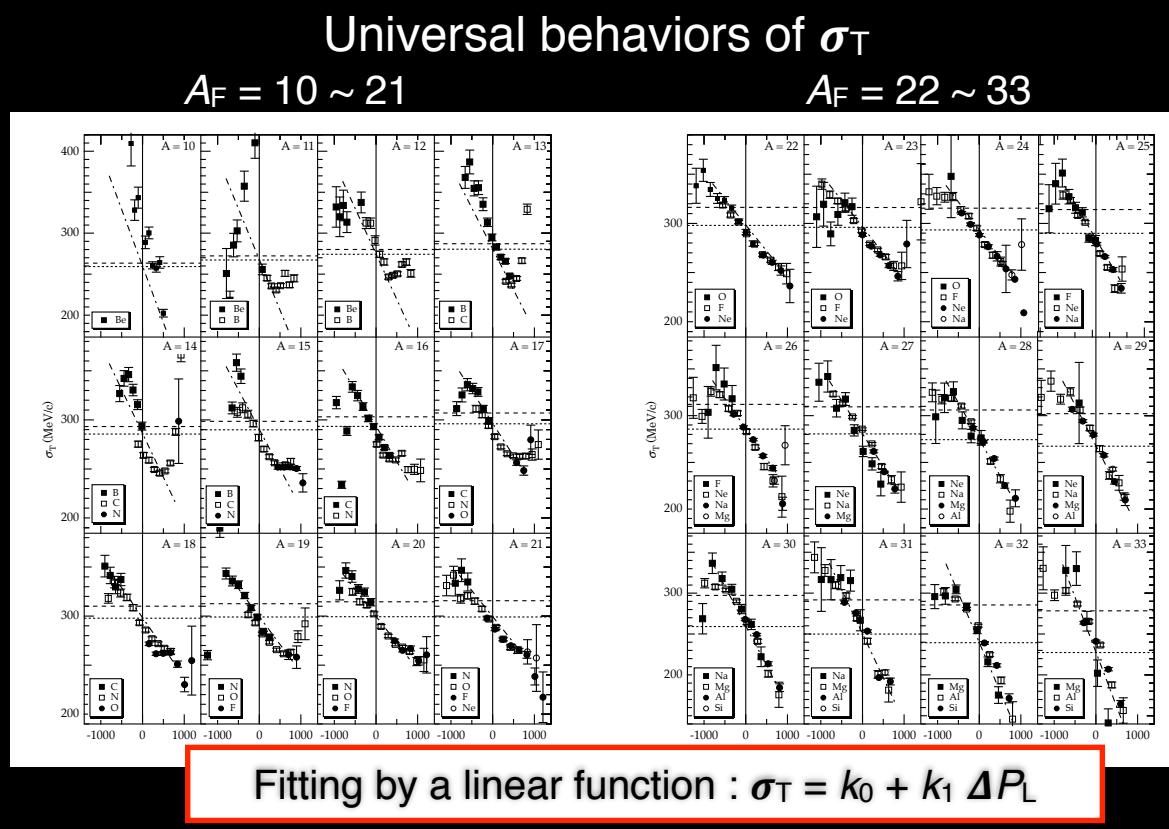
Agreement with reference values

GH@primary beam velocity

GH+ Bibber@center of P_L dist.

S. Momota et al., Phys. Rev. C 92, 024608 (2015).

Correlation between σ_T and ΔP_L

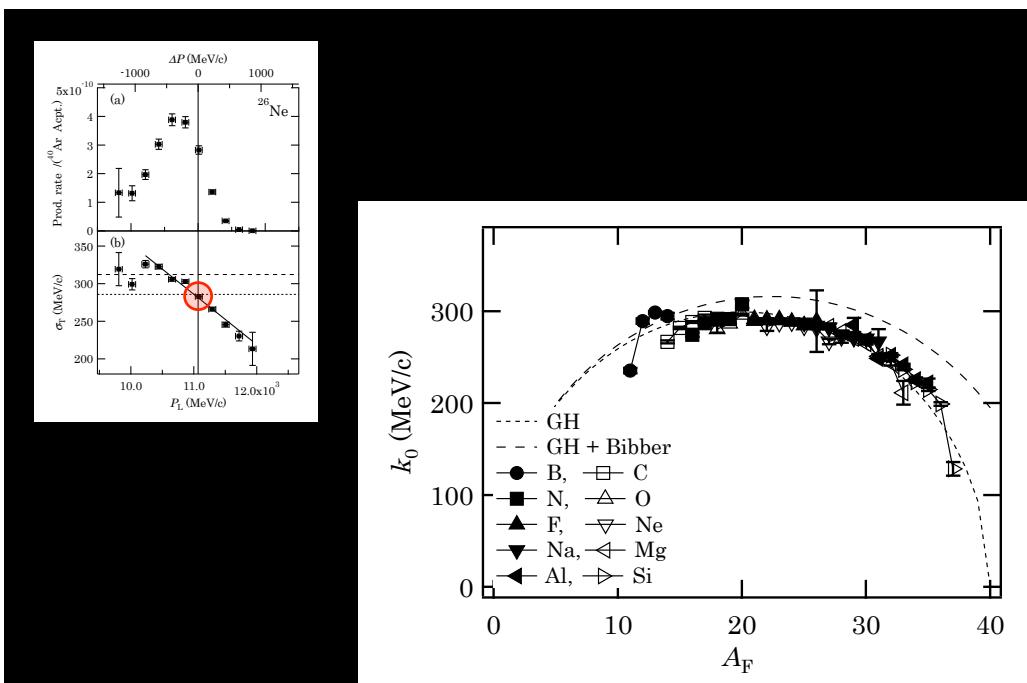


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σ_T at projectile velocity : k_0



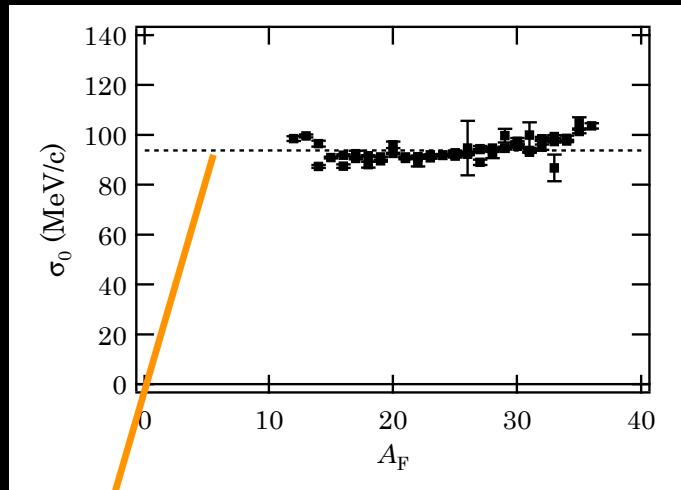
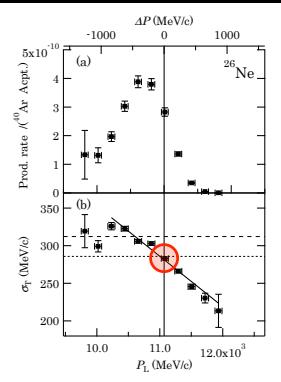
- ✓ Good agreement with Goldhaber formulation
No additional dispersions are not needed.

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Reduced width : σ_0



Av. = $93.6 \pm 1.3 \text{ MeV}/c$

$93.5 \pm 2.6 \text{ MeV}/c$

obtained from P_L dist., $^{40}\text{Ar}(90 \text{ MeV/u}) + \text{Be}$
M. Notani et al., Phys. Rev. C 76 (2007) 044605.

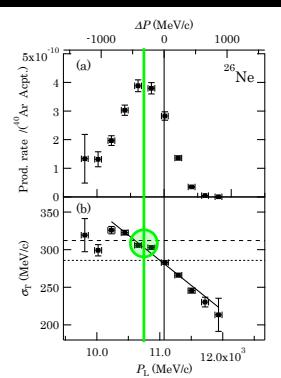
✓ Good agreement with σ_0 obtained from P_L dist.

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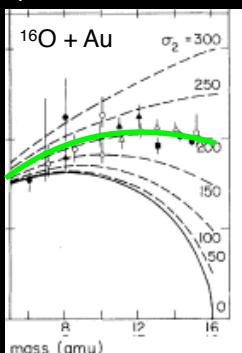
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σ_T at center of P_L distribution

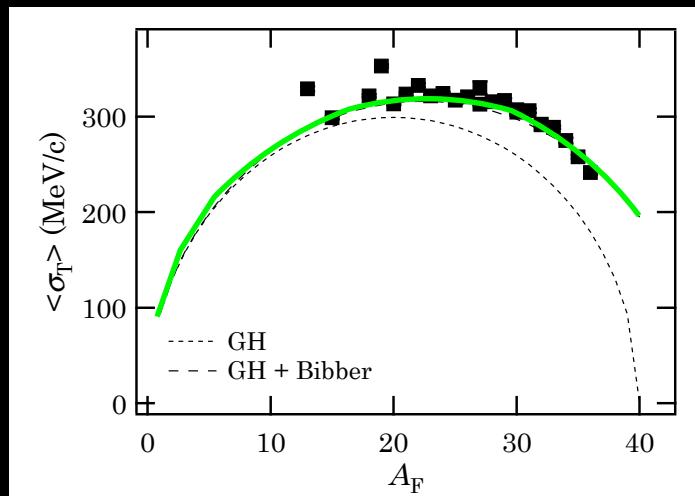


σ_T at $E=92.5 \text{ MeV/u}$



In order to compare with the previous results,
most probable σ_T , $\langle \sigma_T \rangle$, is introduced.

$\langle \sigma_T \rangle$: σ_T at center of P_L distribution



✓ Consistent with previous results on σ_T

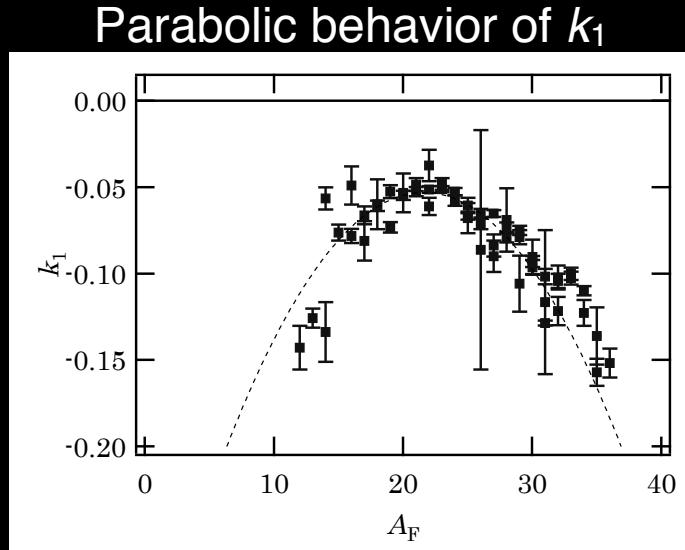
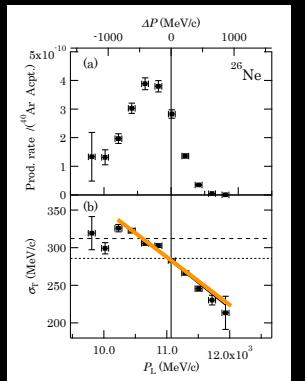
K van Bibber et al.,
Phys. Rev. Lett. 43 (1979) 840.

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Slope parameter : k_1



Fitting by a quadratic function :
 $k_1 = -0.384 + 0.0273A_F + 0.000631A_F^2$

Empirical formulation of σ_T

Width of P_T distribution : σ_T

Monotonically decreasing with velocity

$$\sigma_T = k_0 + k_1 \Delta P_L$$

σ_T at projectile velocity : k_0

$$\sigma_T(\Delta P_L=0) = \sigma_L = \sigma_{GH}$$

Slope parameter : k_1

Depends on A_F

$$k_1 = -0.384 + 0.0273A_F + 0.000631A_F^2$$

Microscopic reaction model
can reproduce behaviors of σ_T ?
can reveal origin of the behaviors?

Microscopic reaction model

Collective features

-> Additional dispersion and deceleration effect

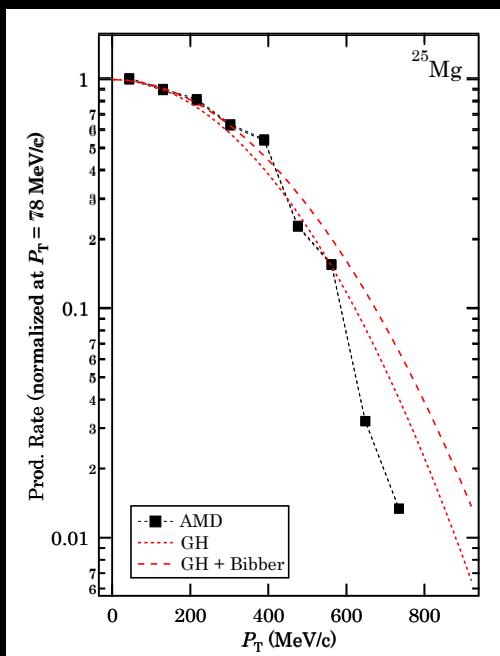


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P_T distribution obtained by simulation

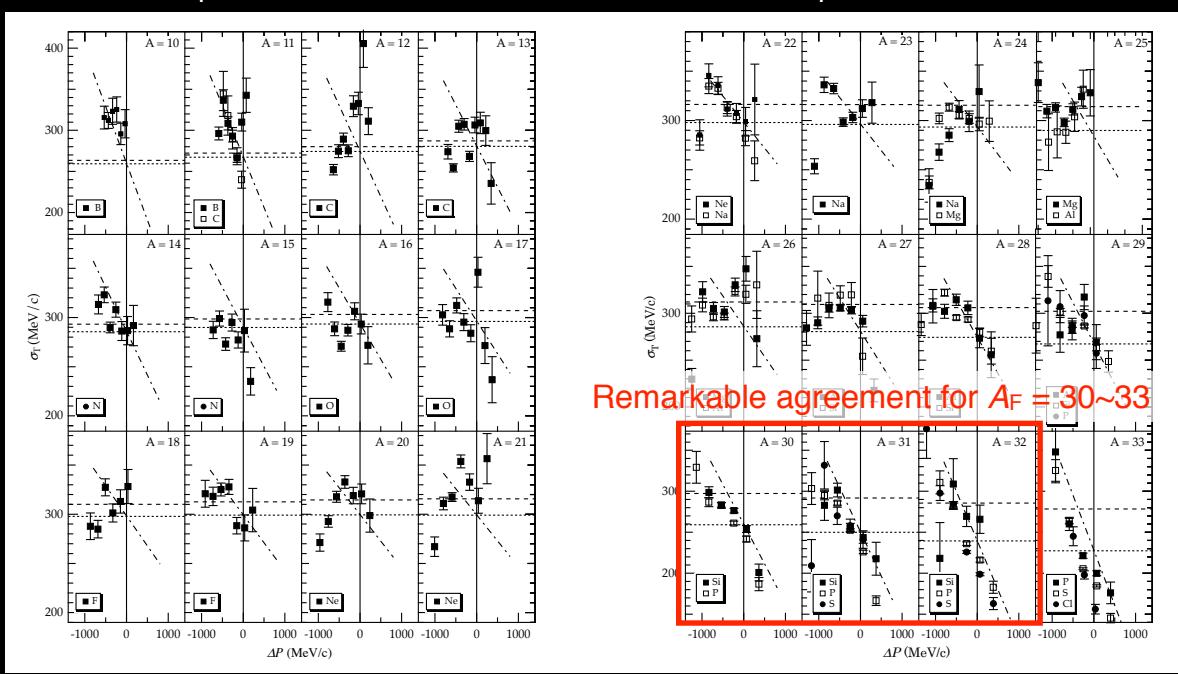


P_T distribution obtained from
AMD + SD calculation

Gaussian-like distribution
Width is consistent with
conventional values

Fitting with a Gaussian function
as for experimental results
→ $\sigma_T(\text{AMD})$

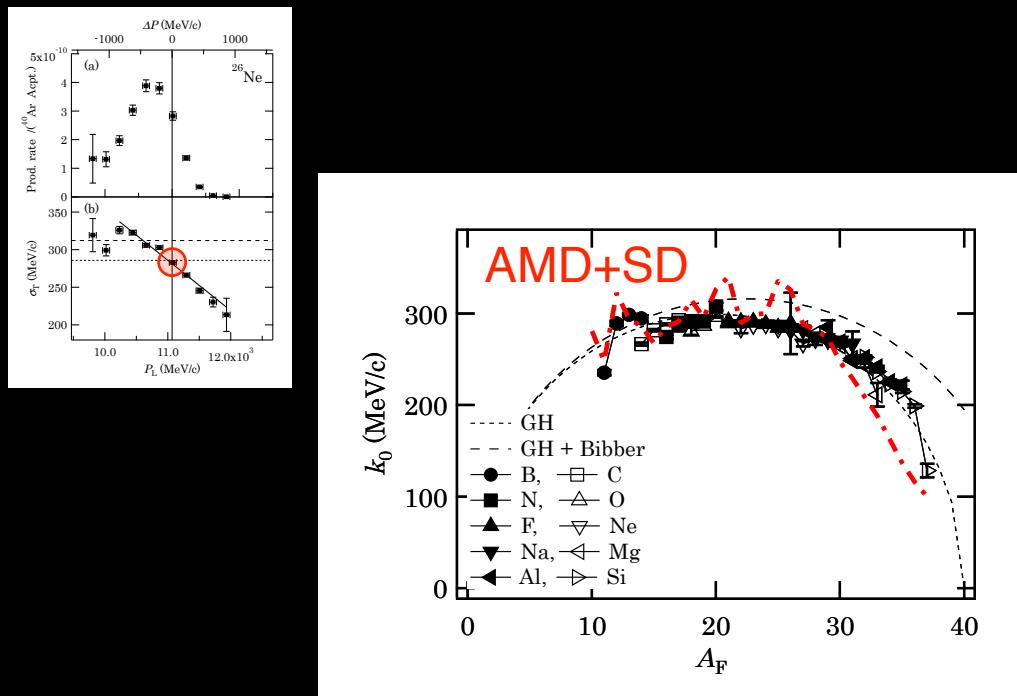
Correlation between σ_T and ΔP_L



AMD calculation roughly reproduces behaviors of σ_T .
Fitting with a linear function $\rightarrow k_0, k_1$

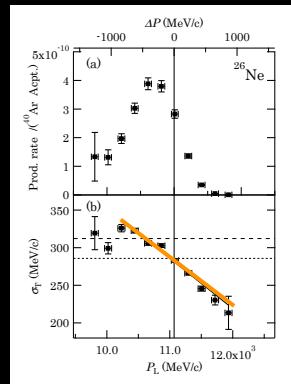
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σ_T at projectile velocity : k_0

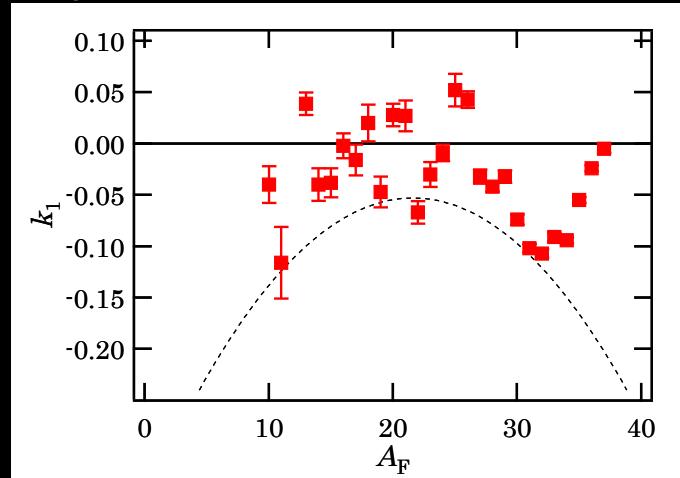


Not so bad,
but systematically underestimate k_0 at $A_F > 30$.

Slope parameter : k_1



Parabolic trend of k_1 obtained from experimental results

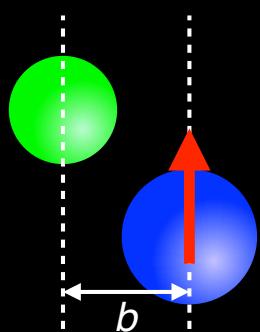


AMD calculation roughly reproduces negative values for k_1 .
Large scattering prevents further investigations on k_1 .

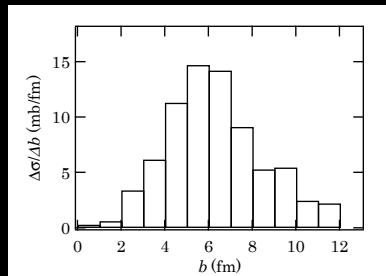
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b -dependent P_L distribution

Impact parameter \leftrightarrow Collective/dissipative nature
Therefore, P_L distribution is expected to depends on impact parameter.



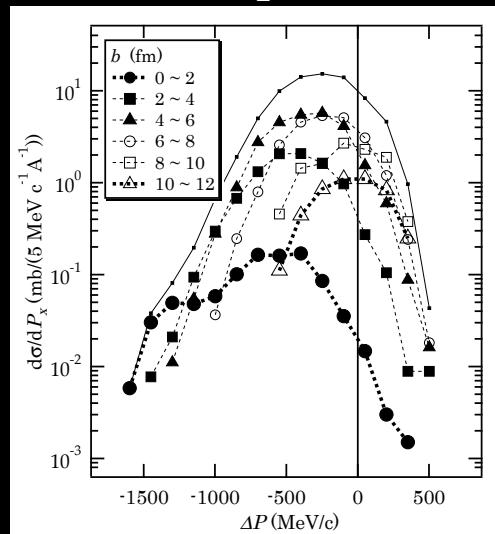
Prod. rate of ^{30}Si vs. b



$$r(^{40}\text{Ar}) + r(^9\text{Be}) = 6.6 \text{ fm}$$

$$r = r_0 A^{1/3}, r_0 = 1.2 \text{ fm}$$

Calculated P_L distribution

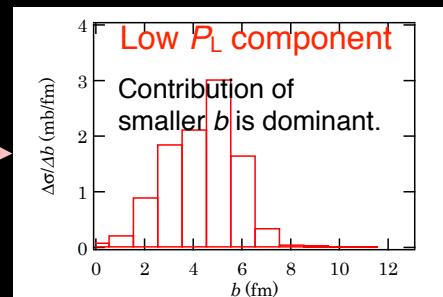
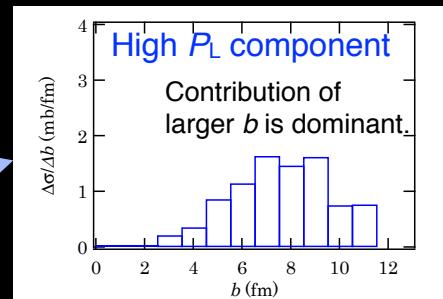
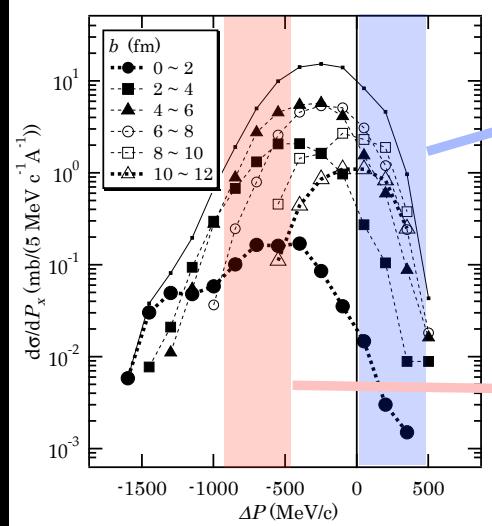


E -dissipation is promoted for small b .

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Contribution of impact parameter

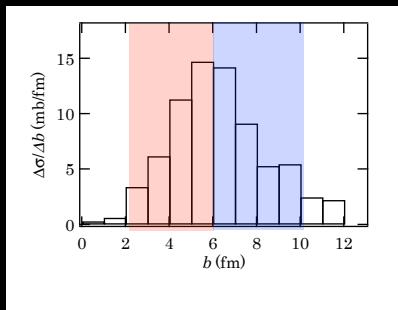
b -dependent P_L distribution



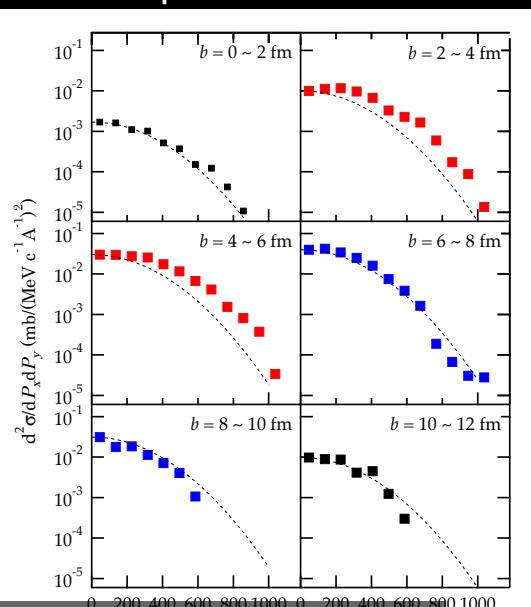
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b -dependent P_T distribution

Prod. rate vs. b



b -dependent P_T dist.



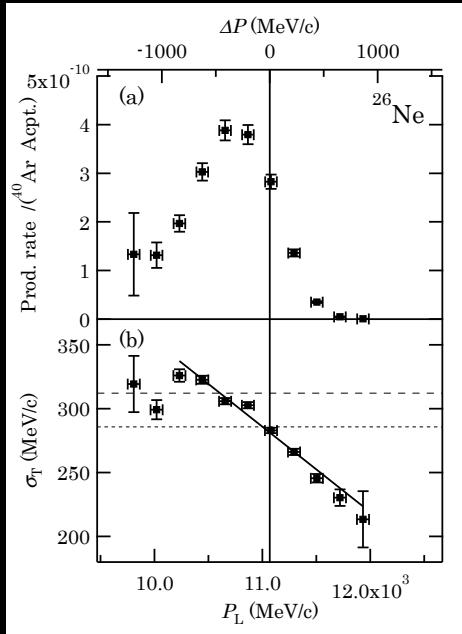
The width of P_T distribution is larger for small b .

→ v -dependent σ_T would be originated from contribution of impact parameters.

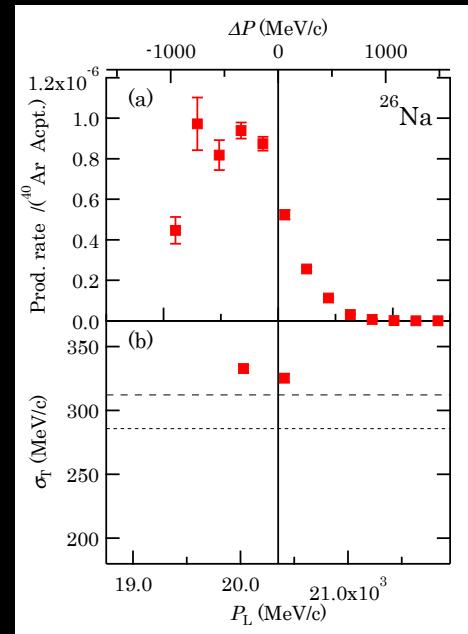
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P_T distribution at higher energy

$E = 95 \text{ MeV/u}$
 $^{40}\text{Ar} + \text{Be} \rightarrow ^{26}\text{Ne} + X$



$E = 290 \text{ MeV/u}$
 $^{40}\text{Ar} + \text{Al} \rightarrow ^{26}\text{Na} + X$

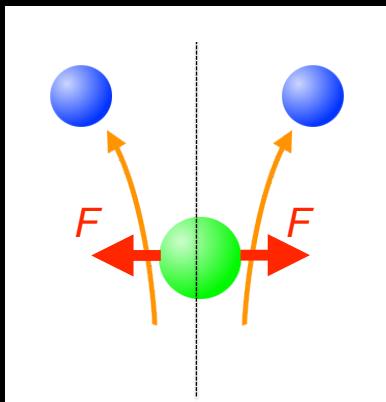


Velocity dependence is not remarkable.

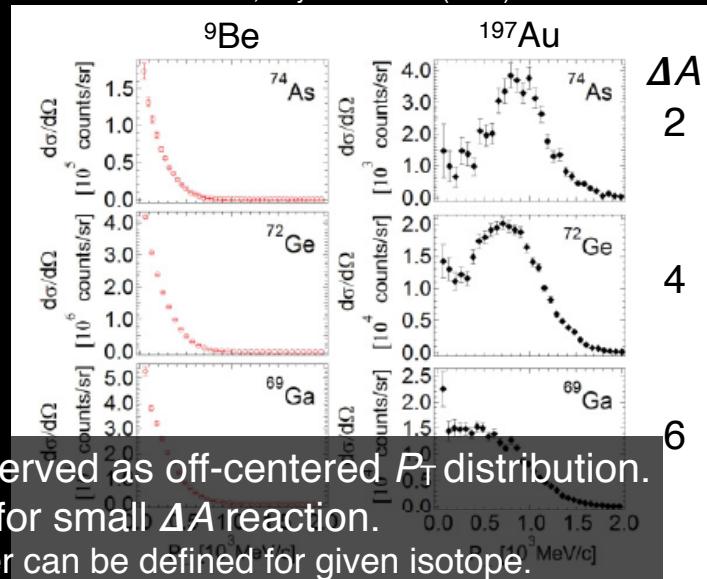
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P_T distribution with heavier target

Dominant contribution of repulsive Coulomb force



P_T distribution of fragments
 $^{76}\text{Ge}(130 \text{ MeV/nucleon}) + ^{9}\text{Be}, ^{197}\text{Au}$
K. Meierbach et al., Phys. Rev. C 85(2012) 034608



- Deflection effect was observed as off-centered P_T distribution.
- Deflection is remarkable for small ΔA reaction.
Specified impact parameter can be defined for given isotope.

Possibility to investigation proximity potential for heavy reaction system

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Conclusions

- Remarkable **correlation** between width of P_T distribution (σ_T) and fragment velocity (ΔP_L) has been observed at $E = 95$ MeV/u.

1. Simple **formulation** : $\sigma_T = k_0 + k_1 \Delta P_L$

2. Comparison with **previous results**

$$\sigma_T(\Delta P_L = 0) = \sigma_{GH}$$

Width of P_L distribution

$$\sigma_T \text{ at center of } P_L \text{ dist.} = \sigma_D$$

Conventionally used value

3. Important contribution of **impact parameter** to understand observed correlation

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Acknowledgements

- **p178n@RIPS collaboration**

RCNP, Japan : I. Tanihata

Tsukuba Univ., Japan : A. Ozawa

The Enrico Fermi Inst., USA : M. Notani

RIKEN, Japan : K. Yoshida, K. Morimoto, A. Yoshida

Saitama Univ., Japan : T. Yamaguchi

Hitachi Co., Japan : T. Onishi

KEK, Japan : Y. X. Watanabe

IMP, China : Z. Liu, and

RIKEN Ring Cyclotron staff and crew

- **AMD calculation**

Tohoku Univ., Japan : A. Ono

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